

# $G_2$ -geometry on solvmanifolds

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(joint work with Ilka Agricola & Simon Chiossi)

## Common sector of Type II string equations

- Strominger (1986):  $(M, g)$   $n$ -dimensional Riemannian spin manifold with a 3-form  $T$  (field strength), a spinor field  $\Psi$  (supersymmetry) and a function  $\Phi$  (dilaton).

- Bosonic equations:

$$R_{ij}^g - \frac{1}{4} T_{imn} T_{jmn} + 2 \nabla_i^{LC} \delta_j \phi = 0, \quad \delta(e^{-2\phi} T) = 0.$$

- Fermionic equations:

$$(\nabla_X^{LC} + \frac{1}{4} i_X T) \cdot \psi = 0, \quad T \cdot \Psi = 2d\phi \cdot \Psi.$$

It is impossible to look for solutions on a fixed manifold.

Geometric meaning of the 3-form  $T$ ?

The first fermionic equation means that  $\psi$  is parallel with respect to a new connection:

$$\nabla_X Y = \nabla_X^{LC} Y + \frac{1}{2} T(X, Y, -).$$

$\Rightarrow T =$  torsion of  $\nabla$  and the equations become:

$$\text{Ric}^\nabla + \frac{1}{2} \delta T + 2 \nabla^{LC} d\phi = 0, \quad \delta(e^{-2\phi} T) = 0$$

$$\nabla \psi = 0, \quad T \cdot \psi = 2d\phi \cdot \psi$$

Remarks:

- Bosonic equations generalize Einstein equations of general relativity.
- Calabi-Yau and Joyce manifolds are exact solutions with  $T = 0$  and  $\phi = \text{const.}$

$\Rightarrow$  If  $\dim M = 7$ , integrable  $G_2$  structures are exact solutions and non-integrable  $G_2$ -structures can be studied using metric connections with skew-symmetric torsion.

- Can the same 7-manifold  $(M, g)$  carry integrable and non-integrable  $G_2$  structures *simultaneously*?

- There are 7-manifolds with a parallel spinor for  $\nabla^{LC}$  admitting a covariantly constant spinor for some other  $\nabla$ ?

Aim The answer is YES to both questions. The manifold will be a rank-one solvable extension of a 6-dimensional nilpotent Lie group ( $\Rightarrow$  non compact).

## Connections with skew-symmetric torsion

$(M^n, g)$  oriented Riemannian manifold

$\nabla$ : metric connection

$\nabla_X Y = \nabla_X^{LC} Y + A(X, Y)$ ,  
with  $A \in \mathbb{R}^n \otimes \Lambda^2(\mathbb{R}^n) \Rightarrow A$  can be identified  
with the torsion  $T$  of  $\nabla$ .

If  $A \in \Lambda^3(\mathbb{R}^n) \Rightarrow$

$$\nabla_X Y = \nabla_X^{LC} Y + \frac{1}{2}T(X, Y, -)$$

and  $\nabla$  is called a metric connection with skew-symmetric torsion.

$\Rightarrow$  The  $\nabla$ -Killing vector fields coincide with the Riemannian Killing vector fields.

## Lifting metric connections to $SM$

$$\nabla_X Y = \nabla_X^{LC} Y + A_X Y,$$

with  $A_X = i_X T \in \mathfrak{so}(n) \cong \Lambda^2(\mathbb{R}^n) \Rightarrow$   
 $A_X = \sum_{i < j} \alpha_{ij} e_i \wedge e_j.$

Since the lift into  $\mathfrak{spin}(n)$  of  $e_i \wedge e_j$  is  $\frac{e_i \cdot e_j}{2}$ ,  
 $A_X$  defines an element in  $\mathfrak{spin}(n)$  ( $\Rightarrow$  an endomorphism on  $SM$ ).

The action of  $A_X$  is:

$$A_X Y = i_Y A_X \text{ on vectors } (A_X \in \mathfrak{so}(n))$$

$$A_X \psi = \frac{1}{2} A_X \cdot \psi \text{ on spinors } (A_X \in \mathfrak{spin}(n))$$

“ $\cdot$ ” is the Clifford product of a  $k$ -form by a spinor

$\Rightarrow$  the lift of  $\nabla$  on  $SM$  is

$$\nabla_X \psi = \nabla_X^{LC} \psi + \frac{1}{2} A_X \cdot \psi.$$

## Non existence theorems

If the dilaton function  $\phi = \text{const}$

- Bosonic equations:  $Ric^\nabla = 0, \delta T = 0$
- Fermionic equations:  $\nabla \psi = 0, T \cdot \psi = 0$

Theorem A full solution of Strominger's model with  $\phi = \text{const}$  satisfies necessarily  $T = 0$  or  $\psi = 0$ .

[Compact case, Agricola.

General case, Agricola, Friedrich, Nagy, Puhle.]

Vanishing Theorem [Agricola-Friedrich]

$(M^n, g, T)$  compact spin with  $Scal^g \leq 0$ . If  $dT$  acts on spinors as a non positive endomorphism and  $\exists \psi \neq 0$  solution of

$$\nabla_X \psi = \nabla_X^{LC} \psi + (i_X T) \cdot \psi = 0$$

then  $T = 0 = Scal^g$  and  $\nabla^{LC} \psi = 0$ .

Corollary On a Calabi-Yau or Joyce manifold a torsion connection such that  $dT = 0$  can have parallel spinor only for  $T = 0$ .

$\Rightarrow$  rigidity of vacuum solutions under deformation of the connection.

The Vanishing Theorem applies also to nilmanifolds  $\Gamma \backslash G$ .

We will give 7-dimensional examples for which the rigidity theorem does not hold.

Starting point There exist (incomplete) metrics of holonomy  $G_2$  with 2-step nilpotent isometry group  $N^6$  acting on orbits of codim 1 [Gibbons–Lü–Pope–Stelle].

Such metrics are ‘scale invariant’ meaning that there is a homothetic Killing vector field  $X$  ( $L_X g = c g$ ) and locally conformally equivalent to homogeneous metrics on rank-one solvable extensions of  $N^6$  [Chiossi, F].

⇒ classification of 6-dimensional nilpotent Lie algebras endowed with an  $SU(3)$ -structure and a “Hermitian” derivation  $D$  such that the  $G_2$ -structure on the solvable extension is conformally parallel (integrable).

## SU(3) structures

Let  $N^6$  be a (real) 6-dimensional manifold.  
An  $SU(3)$  reduction is given by

$$(J, \omega, h)$$

plus a choice of

$$\eta = \eta^+ + i\eta^- \in \Lambda^{3,0}$$

One can choose an orthonormal basis of 1-forms such that

$$\begin{aligned}\omega &= e^{14} - e^{23} + e^{56}, \\ \eta &= (e^1 + ie^4) \wedge (e^2 - ie^3) \wedge (e^5 + ie^6)\end{aligned}$$

with  $\eta^\pm \wedge \omega = 0$ .

Note  $\eta^+$  determines  $J$  and  $\eta^- = J\eta^+$ .

## Intrinsic torsion for $SU(3)$

$$\text{Hol}(N^6, h) \subseteq SU(3) \iff \begin{cases} d\omega = 0, \\ d\eta^+ = 0, \\ d\eta^- = 0. \end{cases}$$

The failure of this condition is measured by the intrinsic torsion (tensor)  $\tau$  belonging to the space

$$\begin{aligned} T^* \otimes \mathfrak{su}(3)^\perp &\cong T^* \otimes (\llbracket \Lambda^{2,0} \rrbracket \oplus \mathbb{R}) \\ &= \mathcal{W}_1 \oplus \mathcal{W}_2 \oplus \mathcal{W}_3 \oplus \mathcal{W}_4 \oplus \mathcal{W}_5 \end{aligned}$$

where  $\mathcal{W}_j$  are the so-called ‘Gray–Hervella classes’ for  $j \leq 4$ .

If  $d\omega^2 = d\eta^+ = 0 \Rightarrow$  the  $SU(3)$  structure is *half-flat*.

## $G_2$ structures

Given  $M^7$  such that

$$T_m^* M^7 = T_n^* N^6 \oplus \mathbb{R}$$

(e.g. a Riemannian product  $M^7 = N^6 \times \mathbb{R}$ )  
and an  $SU(3)$ -structure  $(\omega, \eta^+)$ , define

$$\begin{aligned}\varphi &= \omega \wedge e^7 + \eta^+, \\ * \varphi &= \eta^- \wedge e^7 + \frac{1}{2} \omega^2,\end{aligned}$$

where  $e^7 = dt$  ( $t$  a coordinate on  $\mathbb{R}$ ).

The associated metric  $g$  has an orthonormal basis for which

$$\varphi = e^{147} - e^{237} + e^{567} + e^{125} + e^{136} + e^{246} - e^{345}$$

has isotropy  $G_2 \subseteq SO(7)$ .

The intrinsic torsion of a  $G_2$  structure can be identified with  $\nabla^{LC}\varphi$  and is encoded into the exterior derivatives  $d\varphi$ ,  $d*\varphi$  by

$$\begin{aligned}d\varphi &= \tau_0 * \varphi + 3\tau_1 \wedge \varphi + *\tau_3, \\d*\varphi &= 4\tau_1 \wedge *\varphi + \tau_2 \wedge \varphi\end{aligned}$$

where  $\tau_i$  is an intrinsic torsion  $i$ -form.

$$\begin{aligned}\nabla\varphi \in T^* \otimes \mathfrak{g}_2^\perp &\cong \mathcal{X}_1 \oplus \mathcal{X}_2 \oplus \mathcal{X}_3 \oplus \mathcal{X}_4 \\ &\cong \mathbb{R} \oplus \mathfrak{g}_2 \oplus S_0(\mathbb{R}^7) \oplus \mathbb{R}^7\end{aligned}$$

[Fernández-Gray]

class	type	conditions
$\mathcal{X}_4$	conformally parallel	$\tau_0 = \tau_2 = \tau_3 = 0$
$\mathcal{X}_1 \oplus \mathcal{X}_3$	cocalibrated	$\tau_1 = \tau_2 = 0$
$\mathcal{X}_1 \oplus \mathcal{X}_3 \oplus \mathcal{X}_4$	$G_2T$	$\tau_2 = 0$

Theorem [Friedrich-Ivanov]

$\tau_2 = 0 \iff \exists!$  a torsion connection  $\nabla$  such that  $\nabla\varphi = 0$

$\iff \exists$  a vector field  $X$  such that  $\delta\varphi = -i_X\varphi$ .

The resulting 3-form torsion is

$$T = \frac{7}{6}\tau_0\varphi - *d\varphi + *(4\tau_1 \wedge \varphi)$$

and  $\nabla$  admits (at least) one parallel spinor.

$$T \in \Lambda^3(\mathbb{R}^7) = \mathbb{R} \oplus S_0(\mathbb{R}^7) \oplus \mathbb{R}^7.$$

$(M^7, \varphi)$  is locally conformal parallel iff

$$\begin{aligned}d\varphi &= 3\tau_1 \wedge \varphi, \\d*\varphi &= 4\tau_1 \wedge *\varphi.\end{aligned}$$

In this case,  $T = *(\tau_1 \wedge \varphi)$ .

$G_2$  can be lifted to a subgroup of  $Spin(7)$ .

$\Delta_7$ : complex vector space of all 7-dimensional spinors  $\Rightarrow \Delta_7$  is a complex  $Spin(7)$ -representation and the complexification of a real representation.

$$\Delta_7 \cong \mathbb{R} \oplus \mathbb{R}^7 \text{ (under } G_2\text{)}.$$

We will use the real representation obtained by

$$e_1 = +E_{18} + E_{27} - E_{36} - E_{45},$$

$$e_2 = -E_{17} + E_{28} + E_{35} - E_{46},$$

$$e_3 = -E_{16} + E_{25} - E_{38} + E_{47},$$

$$e_4 = -E_{15} - E_{26} - E_{37} - E_{48},$$

$$e_5 = -E_{13} - E_{24} + E_{57} + E_{68},$$

$$e_6 = +E_{14} - E_{23} - E_{58} + E_{67},$$

$$e_7 = +E_{12} - E_{34} - E_{56} + E_{78},$$

$E_{ij}$  standard basis of  $\mathfrak{so}(7)$ .

## Solvable Extensions

Let  $(\mathfrak{n}, [\cdot, \cdot]_{\mathfrak{n}}, \langle \cdot, \cdot \rangle')$  be a metric nilpotent Lie algebra.

Then  $(\mathfrak{s} = \mathfrak{n} \oplus \mathfrak{a}, \langle \cdot, \cdot \rangle)$  is a metric solvable extension of  $(\mathfrak{n}, \langle \cdot, \cdot \rangle')$  if  $[\cdot, \cdot]$  restricted to  $\mathfrak{n}$  coincides with  $[\cdot, \cdot]_{\mathfrak{n}}$  and  $\langle \cdot, \cdot \rangle|_{\mathfrak{n} \times \mathfrak{n}} = \langle \cdot, \cdot \rangle'$ .

One says that  $\mathfrak{s}$  is 'standard' if  $\mathfrak{a} = (\mathfrak{s}^1)^\perp$  is abelian, and  $\dim \mathfrak{a}$  is called the rank.

If the rank is 1 with  $\mathfrak{a} = \langle A \rangle$ , the extension is of *Iwasawa type* if

- (i)  $\text{ad}_A \neq 0$  is self-adjoint with respect to  $\langle \cdot, \cdot \rangle$ ,
- (ii)  $\text{ad}_A|_{\mathfrak{n}}$  is positive-definite.

The study of standard solvable Lie algebras with Einstein metrics reduces to rank-one metric solvable extensions

$$(\mathfrak{g} = \mathfrak{n} \oplus \mathbb{R}H, \langle \cdot, \cdot \rangle)$$

with  $\langle H, \mathfrak{n} \rangle = 0$ ,  $\|H\| = 1$  and Lie bracket

$$\begin{cases} [H, X] = DX, \\ [X, Y] = [X, Y]_{\mathfrak{n}} \end{cases}$$

for some  $\mathfrak{n}$  and  $D \in \text{Der}(\mathfrak{n})$  [Heber].

All nilpotent Lie groups of  $\dim \leq 6$  admit a rank-one Einstein solvable extension [Lauret, Will].

## Classification theorem [Chiossi–F]

Let  $N$  be a nilpotent Lie group of dimension 6 with an invariant  $SU(3)$  structure  $(\omega, \eta^+)$ . Suppose that there is a derivation  $D$  of the Lie algebra  $\mathfrak{n}$  such that  $(DJ)^2 = (JD)^2$ .

Then on the solvable extension

$$\mathfrak{s} = \mathfrak{n} \oplus \langle e_7 \rangle$$

with  $\text{ad}_{e_7} = D$ , the  $G_2$ -structure

$$\varphi = \omega \wedge e^7 + \eta^+$$

is conformally parallel iff  $\mathfrak{n}$  is

- (i) either  $\mathbb{R}^6$  or 2-step nilpotent, and
- (ii) not isomorphic to  $\mathfrak{h}_3 \oplus \mathfrak{h}_3$ .

## Sketch of the proof

• Since the derivation  $D = \text{ad}_{e_7}$  preserves the orthogonal splitting  $\mathfrak{n}^1 \oplus (\mathfrak{n}^1)^\perp$ , one can suppose that  $D$  is diagonalizable by a unitary basis  $\{e_1, \dots, e_6\}$  with  $e_1 \in (\mathfrak{n}^1)^\perp$ . The Maurer-Cartan equations of  $\mathfrak{s}$  are

$$\begin{cases} de^j = \hat{d}e^j + C_j e^{j7}, & 1 \leq j \leq 6 \\ de^7 = 0 \end{cases}$$

where  $\hat{d} = d|_{\Lambda^*\mathbb{R}^6}$ ,  $\hat{d}e^1 = 0$ , and  $C_j \neq 0$ .

•  $\varphi$  is conformally parallel iff  $\exists m \neq 0$  such that

$$\begin{aligned} d\omega \wedge e^7 + d\eta^+ &= -3m\eta^+ \wedge e^7, \\ d\eta^- \wedge e^7 + \omega \wedge d\omega &= 2m\omega^2 \wedge e^7. \end{aligned} \quad (*)$$

In this case,  $(\omega, \eta^+)$  is necessarily half-flat.

## Consequences

If  $\mathfrak{n}$  is not abelian  $\Rightarrow$

- $J$  cannot be integrable
- $d\omega \neq 0$ .

[Conti-Tomassini studied symplectic half-flat structures on 6-dimensional nilmanifolds  $N$ ]

$dT \neq 0$  unless  $m = 0$ .

[Chiossi–Swann proved that if  $J$  is integrable and  $N \times S^1$  is  $G_2T$ , then  $N$  is balanced]

Corollary Any solvmanifold corresponding to a 6-dimensional 2-step nilpotent Lie algebra, except  $\mathfrak{h}_3 \oplus \mathfrak{h}_3$ , admits a pair of Einstein metrics (one homogeneous and one Ricci-flat).

nilpotent Lie algebra	eigenvalues of $ad_{e_7}$
$(0, 0, e^{15}, 0, 0, 0)$	$-\left(\frac{2m}{3}, m, \frac{4m}{3}, m, \frac{2m}{3}, m\right)$
$(0, 0, e^{15}, e^{25}, 0, e^{12})$	$-\left(\frac{3m}{5}, \frac{3m}{5}, \frac{6m}{5}, \frac{6m}{5}, \frac{3m}{5}, \frac{6m}{5}\right)$
$(0, 0, e^{15} - e^{46}, 0, 0, 0)$	$-\left(\frac{3m}{4}, m, \frac{3m}{2}, \frac{3m}{4}, \frac{3m}{4}, \frac{3m}{4}\right)$
$(0, e^{45}, -e^{15} - e^{46}, 0, 0, 0)$	$-\left(\frac{4m}{5}, \frac{6m}{5}, \frac{7m}{5}, \frac{3m}{5}, \frac{3m}{5}, \frac{4m}{5}\right)$
$(0, e^{45}, e^{46}, 0, 0, 0)$	$-\left(m, \frac{5m}{4}, \frac{5m}{4}, \frac{m}{2}, \frac{3m}{4}, \frac{3m}{4}\right)$
$(0, e^{16} + e^{45}, e^{15} - e^{46}, 0, 0, 0)$	$-\left(\frac{2m}{3}, \frac{4m}{3}, \frac{4m}{3}, \frac{2m}{3}, \frac{2m}{3}, \frac{2m}{3}\right)$

## Corresponding Ricci-flat metrics

If  $(S, \varphi, g)$  is conformally parallel, the change  $\tilde{g} = e^{2f}g$  with  $df = me^7$  produces a Ricci-flat metric in terms of coordinates  $(x_1, \dots, x_6, t)$  where  $e^7 = dt$ .

$\tilde{g}$  is Ricci-flat and admits LC-parallel spinors.

$\varphi$  defines a LC-parallel spinor  $\Psi$  by

$$\varphi(X, Y, Z) = \frac{1}{4} \langle X \cdot Y \cdot Z \cdot \Psi, \Psi \rangle,$$

where “ $\cdot$ ” is the Clifford multiplication and  $\langle, \rangle$  is the scalar product on the spinor bundle.

In terms of the spin representation  $\Delta_7$  one has  $\Psi = (0, 0, 0, 0, 1, 1, -1, 1)$ .

Comparing the metrics with the ones found by GLPS

nilpotent Lie algebra	Hol	LC-parallel spinors
$(0, 0, e^{15}, e^{25}, 0, e^{12})$	$G_2$	1
$(0, e^{45}, -e^{15} - e^{46}, 0, 0, 0)$	$G_2$	1
$(0, e^{16} + e^{45}, e^{15} - e^{46}, 0, 0, 0, )$	$G_2$	1
$(0, 0, e^{15} - e^{46}, 0, 0, 0)$	$SU(3)$	2
$(0, e^{45}, e^{46}, 0, 0, 0)$	$SU(3)$	2
$(0, 0, e^{15}, 0, 0, 0)$	$SU(2)$	4

E.g. The metric corresponding to  $(0, 0, e^{15}, e^{25}, 0, e^{12})$  is locally isometric to the scale-invariant metric on  $M^6 \times \mathbb{R}$ , where

$$M^6 \xrightarrow{T^3} T^3.$$

Example The metric

$$\begin{aligned}\tilde{g} = & e^{-\frac{2}{5}mt}(dx_1^2 + dx_6^2) + e^{-\frac{4}{5}mt}(dx_4^2 + dx_5^2) \\ & + \frac{9}{25}m^2e^{\frac{4}{5}mt}(dx_3 - \frac{2}{3}x_1dx_5 + \frac{2}{3}x_4dx_6)^2 \\ & + \frac{9}{25}m^2e^{\frac{2}{5}mt}(dx_2 + \frac{2}{3}x_4dx_5)^2 + e^{-2mt}dt^2\end{aligned}$$

on the solvmanifold corresponding to

$$(0, e^{45}, -e^{15} - e^{46}, 0, 0, 0)$$

is new!

It is scale invariant with symmetry generated by the homothetic Killing field

$$\begin{aligned}Z = & -\frac{5}{m}\frac{\partial}{\partial t} + 4x_1\frac{\partial}{\partial x_1} + 4x_6\frac{\partial}{\partial x_6} + 3x_4\frac{\partial}{\partial x_4} + 3x_5\frac{\partial}{\partial x_5} \\ & + \frac{21}{5}mx_3\frac{\partial}{\partial x_3} + \frac{18}{5}mx_2\frac{\partial}{\partial x_2}.\end{aligned}$$

## Evolution of $SU(3)$ -structures

If the  $G_2$ -structure

$$\varphi = \omega \wedge e^7 + \eta^+$$

on  $\mathfrak{s}$  is conformally parallel,  $(\omega, \eta^+)$  is half-flat.

Theorem [Hitchin]

If a compact manifold  $N^6$  has a half-flat  $SU(3)$ -structure,  $\exists$  a metric with  $\text{Hol} \subseteq G_2$  on  $N^6 \times I$  for some interval  $I$ .

Proposition [Chiossi-F]

Any of the Ricci-flat metrics  $\tilde{g}$  on the solvable Lie group  $S$  can be obtained evolving the  $SU(3)$  structure on the 2-step nilmanifold  $\Gamma \backslash N^6$ .

## Questions related to spinors:

- Is  $\tilde{g}$  induced by another (not integrable)  $G_2$ -structure?
- Which solvable Lie groups  $(S, \tilde{g})$  admit a parallel spinor for another torsion connection?

We will assume  $T$  has the form

$$T \in \Lambda_{11}^3 = \langle \text{simple forms appearing in } \eta^+, \eta^-, \omega \wedge e^7 \rangle.$$

The Clifford multiplication by  $i_{e_i}T$  has, as an endomorphism, the block structure

$$\begin{pmatrix} 0 & * \\ * & 0 \end{pmatrix}.$$

The Levi-Civita connection  $\tilde{\nabla}^{LC}$  is given by:

$$\tilde{\nabla}_X^{LC} Y = \nabla_X^{LC} Y + df(X)Y + df(Y)X - g(X, Y)\text{grad}f,$$

with  $df = me^7 \Rightarrow \tilde{\nabla}_{e_7}^{LC} X = 0, \forall X \in \mathfrak{n} \Rightarrow$

$$\text{Ker}(\tilde{\nabla}_{e_7}^{LC} + i_{e_7}T) = \text{Ker}(i_{e_7}T).$$

## Parallel Spinors

Reduction Theorem [Agricola, Chiossi, F]

A non-trivial element in the kernel of  $i_{e_7}T$  for  $T \neq 0 \in \Lambda_{11}^3(Y)$  is a linear combination of upper-block forms:

- $\Psi = (a, b, c, d, 0, 0, 0, 0)$  with

$$T_{567} = 0, \quad T_{147} = -T_{237};$$

- $\Psi = (a, b, \mp a, \pm b, 0, 0, 0, 0)$  and

$$T_{147} = -T_{237} + \epsilon T_{567}$$

or lower-block forms:

- $\Psi = (0, 0, 0, 0, e, f, g, h)$  with

$$T_{567} = 0, \quad T_{147} = T_{237}.$$

- $\Psi = (0, 0, 0, 0, e, f, \pm e, \mp f)$  and

$$T_{147} = T_{237} + \epsilon T_{567}.$$

$(S, \tilde{g})$ : solvmanifold corresponding to  
 $(0, 0, e^{15}, e^{25}, 0, e^{12})$ .

$\Psi = (0, 0, 0, 0, 1, 1, -1, 1)$ : unique  $LC$ -parallel spinor

Theorem [Agricola, Chiossi, F]

The equation  $\nabla_X \Psi = 0$  admits precisely 7 solutions for some  $T \neq 0$ , namely:

a) An  $\mathbb{RP}^1$ -parameter family of pairs  $(T_{r,s}, \Psi_{r,s})$  such that  $\nabla^{r,s} \Psi_{r,s} = 0$ ;

for  $r = s$ , one has  $T_{r,s} = 0$  and  $\Psi_{r,s}$  is a multiple of  $\Psi$ .

b) Six isolated solutions in pairs  $(T_i^\epsilon, \Psi_i^\epsilon)$  for  $i = 1, 2, 3$  and  $\epsilon = \pm$ .

The  $G_2$  structures admit precisely one parallel spinor, and are of general type

$\mathbb{R} \oplus S_0(\mathbb{R}^7) \oplus \mathbb{R}^7$ , except for

$r = s$  ( $\varphi_{r,r}$  is parallel) and  $r = -s$  (no  $\mathbb{R}$ -part).

$(S, \tilde{g})$ : solvmanifold corresponding to  
 $(0, e^{45}, -e^{15} - e^{46}, 0, 0, 0)$ .

Theorem [Agricola, Chiossi, F]

If  $\exists \psi \neq 0$  solution of  $\nabla\psi = 0$ , then:

(a)  $\psi$  is a multiple of

$(1 + 2i\epsilon\sqrt{2}, 3, 1 + 2i\epsilon\sqrt{2}, -3, 0, 0, 0, 0)$  and

$T = \frac{2}{3}[-2e^{126} + e^{135} - 4e^{245} + e^{346}] + i\epsilon\sqrt{2}[e^{125} + e^{136} + e^{246} + e^{345}] + \frac{2}{3}i\epsilon\sqrt{2}[-e^{147} - e^{567} + 2e^{237}]$  or

(b)  $\psi$  is a multiple of

$(3, -1 + 2i\epsilon\sqrt{2}, -3, -1 + 2i\epsilon\sqrt{2}, 0, 0, 0, 0)$  and

$T = \frac{2}{3}[e_{126} - e_{135} + 4e_{245} - 2e_{346}] + i\epsilon\sqrt{2}[-e_{125} + e_{136} + e_{246} - e_{345}] + \frac{2}{3}i\epsilon\sqrt{2}[-e_{147} + e_{567} + 2e_{237}]$  or

(c)  $\psi$  is a multiple of

$(0, 0, 0, 0, 1 + 2i\epsilon\sqrt{2}, 3, 1 + 2i\epsilon\sqrt{2}, -3)$  and

$T = \frac{2}{3}[e_{126} - 2e_{135} + 4e_{245} - e_{346}] + i\epsilon\sqrt{2}[e_{125} + e_{136} - e_{246} - e_{345}] + \frac{2}{3}i\epsilon\sqrt{2}[e_{147} - e_{567} + 2e_{237}]$ .

Question Any physical interpretation of the complex solutions?